

# The 1D wave equation.

○  $U_{tt} = c^2 U_{xx}$

• Recall that we need two initial conditions and two B.C.'s.

Eg.

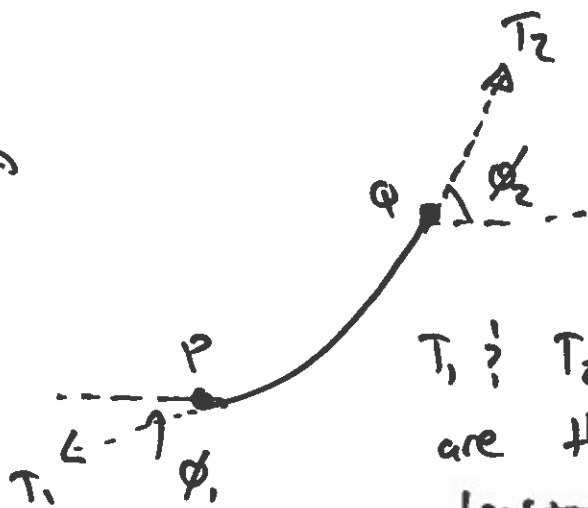
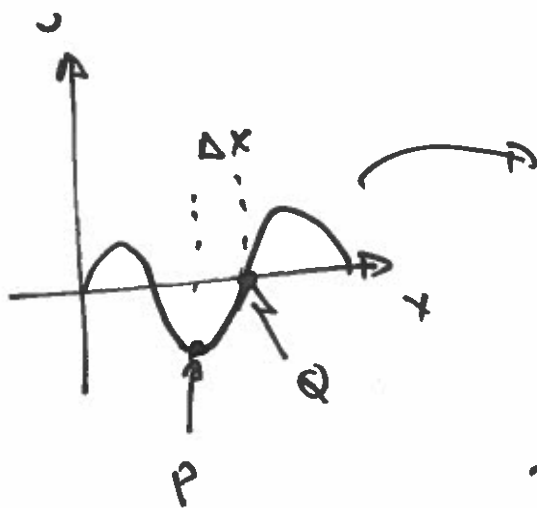
$$\left. \begin{aligned} U(x, 0) &= f(x) \\ U_t(x, 0) &= g(x) \end{aligned} \right\} \text{I.C.}$$

$$\left. \begin{aligned} U(0, t) &= 0 \\ U(L, t) &= 0 \end{aligned} \right\} \text{B.C.}$$

○  $U(L, t) = 0$

## Physical model.

1) Vibrating string.



$T_1, T_2$  are the  ~~tensions~~ tensions at point P & Q.

No motion in the horizontal direction.

$$T_1 \cos \phi_1 = T_2 \cos \phi_2 = T$$

Newton's second law.

$$T_2 \sin \phi_2 - T_1 \sin \phi_1 = \rho \Delta x \frac{\partial^2 u}{\partial t^2}$$

$$e = \frac{kg}{m}$$

(ignoring in plane length/height).

Re-arranging, we get.

$$\tan \phi_2 - \tan \phi_1 = \frac{\rho \Delta x}{T} \frac{\partial^2 u}{\partial t^2}$$

$$\tan \phi_2 = \left. \frac{\partial u}{\partial x} \right|_x$$

$$\tan \phi_1 = \left. \frac{\partial u}{\partial x} \right|_{x+\Delta x}$$

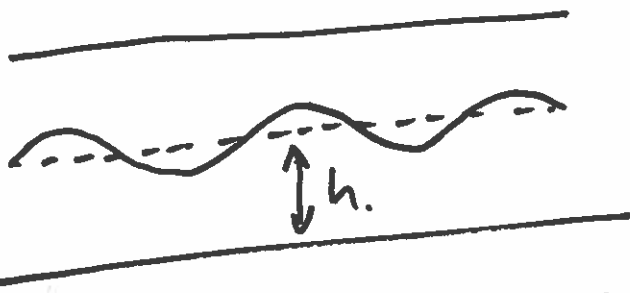
$$\frac{1}{\Delta x} \left( \left. \frac{\partial u}{\partial x} \right|_{x+\Delta x} - \left. \frac{\partial u}{\partial x} \right|_x \right) = \frac{\rho}{T} \frac{\partial^2 u}{\partial t^2}$$

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} \quad c = \sqrt{T/\rho}$$

# Shallow water waves

$$h \sim 10^3 \text{ m.}$$

$$g \sim 10 \text{ m/s}^2$$



Here  $c \approx \sqrt{10^4} = 100 \text{ m/s.}$

The 1D wave equation can be decomposed into a right and left ~~travelling~~ moving wave:

$$\frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = \left[ \frac{\partial}{\partial t} - c \frac{\partial}{\partial x} \right] \left[ \frac{\partial}{\partial t} + c \frac{\partial}{\partial x} \right] u = 0$$

Remember that we can look for an exponential solution:

$$u(x, t) = e^{ikx + i\omega t}$$

Differentiate and sub into wave equation

$$\sigma^2 e^{ikx + \sigma t} = -k^2 c^2 e^{ikx + \sigma t}$$

$$\sigma^2 = -k^2 c^2 \quad \text{or} \quad \sigma = \pm i k c \quad \uparrow \text{dispersion relation}$$

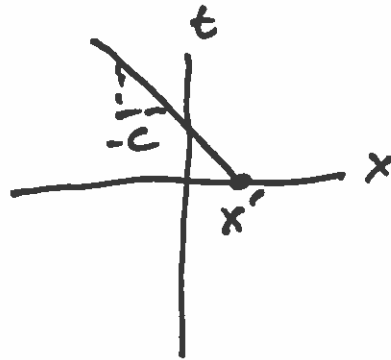
Our guess solution will have the form.

$$u(x, t) = e^{ik(x \pm ct)}$$

### Galilean transform.

• By transforming  $x' = x \pm ct$ , this transforms the frame of reference moving to the left, or to the right at speed  $c$ .

• If the wave is traveling to the left  $x' = x + ct$ .



•  $u(x, t) = G(x+ct)$ . ← left moving wave.

$$u_t = cG' \quad u_x = G'$$
$$u_{tt} = c^2G'' \quad u_{xx} = G''$$

Hence, we get

$$u_{tt} - c^2u_{xx} = c^2G'' - c^2G'' = 0.$$

D'Alembert solution.

$$u_{tt} = c^2u_{xx}.$$

Subject to:  $u(x, 0) = f(x)$ ,  $u_t(x, 0) = g(x)$ .

• Because the wave equation is linear, the general solution is a linear combination of the left and right moving ~~waves~~ waves.

$$u(x, t) = F(x-ct) + G(x+ct).$$

● Apply the IC's.

$$\star f(x) = u(x, 0) = F(x) + G(x).$$

$$g(x) = u_t(x, 0) = -cF'(x) + cG'(x).$$

$$\rightarrow \int_0^x -cF'(x) + cG'(x) dx = \int_0^x g(s) ds + A.$$

$$\hookrightarrow -cF(x) + cG(x) = \int_0^x g(s) ds + A. \quad (1)$$

●  $\star \rightarrow cF(x) + cG(x) = cf(x). \quad (2)$

$$(1) + (2) \rightarrow G(x) = \frac{1}{2} f(x) + \frac{1}{2c} \int_0^x g(s) ds + \frac{A}{2c}$$

$$(2) - (1) \rightarrow F(x) = \frac{1}{2} f(x) - \frac{1}{2c} \int_0^x g(s) ds - \frac{A}{2c}.$$

∴ The & time dependent solution is.

$$u(x, t) = F(x - ct) + G(x + ct).$$

$$u(x,t) = \frac{1}{2} \left\{ f(x-ct) + f(x+ct) \right\} + \frac{1}{2c} \left\{ \int_{x-ct}^x g(s) ds + \int_0^{x+ct} g(s) ds \right\} + \frac{A}{2c} - \frac{A}{2c}.$$

or

$$u(x,t) = \frac{1}{2} \left\{ f(x-ct) + f(x+ct) \right\} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(s) \cdot ds.$$

Left / right waves from  
initial displacement.

depends on  
initial velocity.

Solving the wave equation using separation of variables.

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- Consider the vibrating ~~string~~ string problem

$$v_{tt} = c^2 v_{xx} \quad 0 < x < L.$$

$$\text{B.C: } v(0, t) = v(L, t) = 0.$$

$$\text{I.C: } v(x, 0) = f(x); \quad v_t(x, 0) = g(x).$$

$$v(x, t) = X(x) \cdot T(t).$$

Subing in, we get

$$X(x) \cdot \ddot{T}(t) = c^2 X''(x) \cdot T(t)$$

$$\frac{\ddot{T}(t)}{c^2 T(t)} = \frac{X''(x)}{X(x)} = -\nu^2 \neq 0.$$



$$X] \quad X'' + \mu^2 X = 0.$$

$$\rightarrow X(x) = A \cos(\mu x) + B \sin(\mu x).$$

$$X(0) = X(L) = 0.$$

$$X(0) = 0 = A.$$

$$X(L) = 0 \rightarrow B \sin(\mu L) = 0$$

$$\mu_n = \frac{n\pi}{L} \quad X_n = \sin\left(\frac{n\pi x}{L}\right)$$

$$T] \quad \ddot{T} + \mu^2 c^2 T = 0.$$

$$T(x,t) = A_n \cos(\mu_n c t) + B_n \sin(\mu_n c t)$$

Our solution is then:

$$u(x,t) = \sum_{n=1}^{\infty} \left\{ A_n \cos(\mu_n c t) + B_n \sin(\mu_n c t) \right\} \sin(\mu_n x).$$

To apply the I.C., we take the time derivative

$$u_t(x,t) = \sum_{n=1}^{\infty} \left\{ -A_n \mu_n c \sin(\mu_n c t) + B_n \mu_n c \cos(\mu_n c t) \right\} \cdot \sin(\mu_n x).$$

I.C # 1

$$u(x, 0) = f(x) = \sum_{n=1}^{\infty} A_n \sin(\mu_n x).$$

$$A_n = \frac{2}{L} \int_0^L f(x) \cdot \sin(\mu_n x) dx = b_n^f$$

I.C # 2

$$u_t(x, 0) = \sum_{n=1}^{\infty} B_n \mu_n c \cdot \sin(\mu_n x) = g(x).$$

$$B_n = \frac{2}{\mu_n c L} \int_0^L g(x) \cdot \sin(\mu_n x) dx.$$

$$= \frac{b_n^g}{\mu_n c}$$

So, the solution is:

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ b_n^f \cos(\mu_n c t) + \frac{b_n^g}{\mu_n c} \sin(\mu_n c t) \right\} \sin(\mu_n x).$$

• How does this compare to the

• D'Alembert's solution.

$$\textcircled{1} \quad \sin(A \pm B) = \sin A \cos B \pm \cos A \sin B.$$

$$\textcircled{2} \quad \cos(A \pm B) = \cos A \cos B \mp \sin A \sin B.$$

$$\textcircled{3} \quad \sin A \cos B = \frac{1}{2} \left\{ \sin(A+B) + \sin(A-B) \right\}$$

$$\textcircled{4} \quad \sin A \sin B = \frac{1}{2} \left\{ \cos(A-B) - \cos(A+B) \right\}.$$

$$u(x, t) = \sum_{n=1}^{\infty} \frac{b_n^f}{2} \left\{ \sin[\mu_n(x+ct)] + \sin[\mu_n(x-ct)] \right\}$$

$$+ \sum_{n=1}^{\infty} \frac{b_n^g}{2\mu_n c} \left\{ \cos[\mu_n(x-ct)] - \cos[\mu_n(x+ct)] \right\}$$

Calling the two sums  $u^f$  and  $u^g$ .

$$u(x, t) = u^f + u^g$$

Also, recall that at  $t=0$ , we found that

the ICB are.

$$f^0(x) = \sum_{n=1}^{\infty} b_n^f \sin(\mu_n x). \quad \leftarrow \text{odd extension of } f \text{ cal.}$$

$$g^0(x) = \sum_{n=1}^{\infty} b_n^g \sin(\mu_n x). \quad \leftarrow \text{odd extension of } g \text{ cal.}$$

So,

$$u^f = \sum_{n=1}^{\infty} \frac{b_n^f}{2} \left[ \sin(\mu_n(x+ct)) + \sin(\mu_n(x-ct)) \right]$$

$$= \frac{1}{2} \left\{ f^0(x+ct) + f^0(x-ct) \right\}$$

From (14),

$$g^0(x) = \sum_{n=1}^{\infty} b_n^g \sin(\mu_n x).$$

$$\int_0^x g^0(s) ds = \sum_{n=1}^{\infty} \int_0^x b_n^g \sin(\mu_n s) ds.$$

$$\int_0^x g(s) ds = \sum_{n=1}^{\infty} \frac{-b_n^g}{\mu_n} \cos(\mu_n x) + A.$$

$$u^g = \sum_{n=1}^{\infty} \frac{b_n^g}{2\mu_n c} \cos(\mu_n(x-ct)) - \sum_{n=1}^{\infty} \frac{b_n^g}{2\mu_n c} \cos(\mu_n(x+ct)).$$

$$= \frac{1}{2c} \int_{x-ct}^0 g^{\circ}(s) ds + \cancel{A} + \frac{1}{2c} \int_0^{x+ct} g^{\circ}(s) ds - \cancel{A}$$

$$= \frac{1}{2c} \int_{x-ct}^{x+ct} g^{\circ}(s) ds.$$

$$u(x,t) = u^f + u^g$$

$$= \sum_{n=1}^{\infty} \left\{ b_n^f \cos(\mu_n ct) + \frac{b_n^g}{\mu_n c} \sin(\mu_n ct) \right\} \sin(\mu_n x)$$

$$= \frac{1}{2} \left\{ f^{\circ}(x+ct) + f^{\circ}(x-ct) \right\} + \frac{1}{2c} \int_{x-ct}^{x+ct} g^{\circ}(s) ds.$$

$f^{\circ}(x)$  ;  $g^{\circ}(x)$  are odd extensions.

## Modes of vibration.

$$k_n = \frac{n\pi}{L}$$

frequency & period.

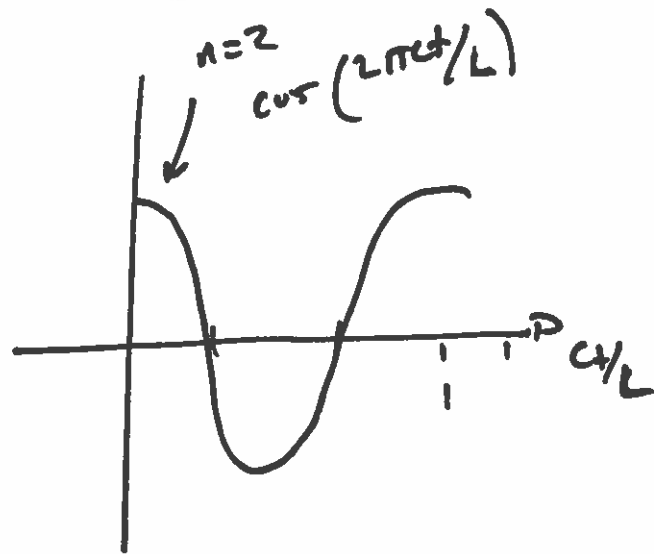
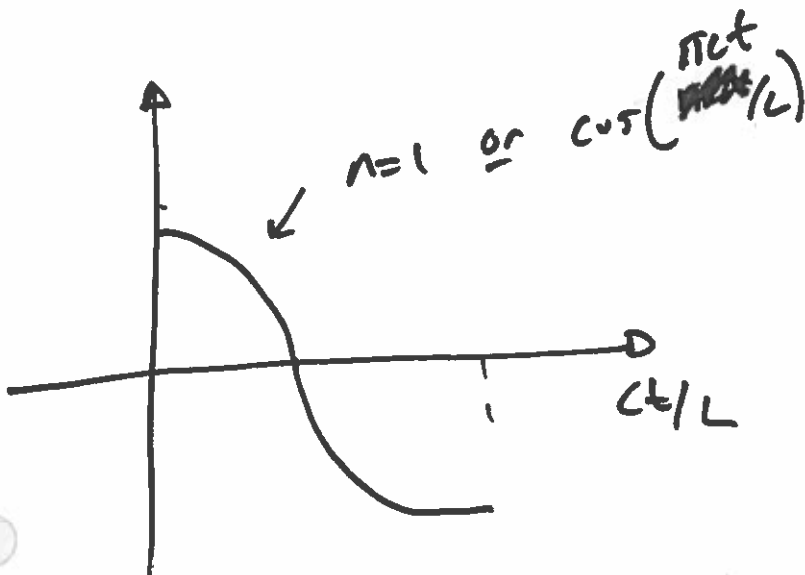
$$\cos[k_n c(t + T_n)] = \cos[k_n ct]$$

$$2\pi = T_n k_n c = 2\pi$$

$$T_n = \frac{2L}{nc} \quad (\text{period})$$

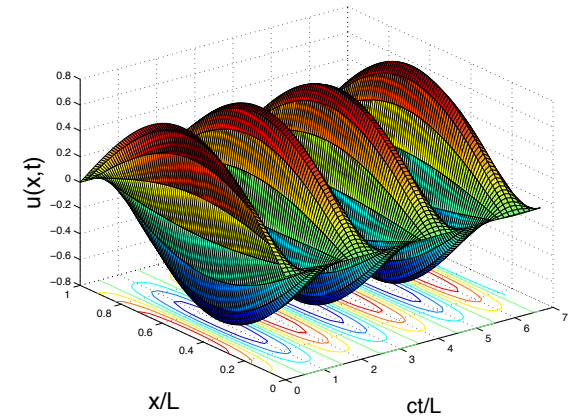
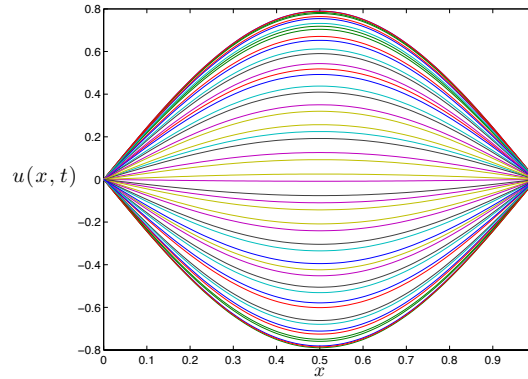
The frequency is:

$$f_n = \frac{1}{T_n} = \frac{nc}{2L}$$

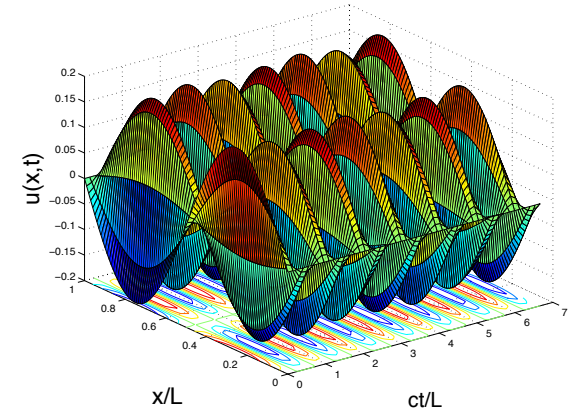
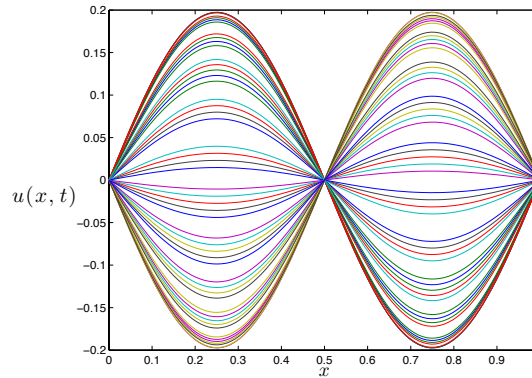


# Lecture 26

Fundamental



First overtone



Second overtone

