



Wave Propagation in Geometrically Reconfigurable Magneto-Elastic Meta-Structures

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Introduction

• Objectives:

- Investigate shape and topology adaptation of periodic lattices
- Investigate effects on:
 - Overall geometry and shape
 - Mechanical properties
 - Wave propagation characteristics
- Approach
 - Exploit bi-stable interactions at the unit cell level to achieve large overall effect at structural level
 - Bi-stability is achieved through magneto-elastic interactions
 - Study is conducted through:
 - Numerical simulations
 - Determination of equilibrium configurations and
 - Bloch analysis of linearized systems

Configurations: reconfigurable hexagonal lattices

Hexagonal Lattice



Re-entrant Lattice



Configurations: Kagome Lattices























Magnetic bi-stable lattices







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Periodic Lattice

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Particle Position

$$egin{aligned} m{r}_{i,nm} &= m{r}_i + nm{d}_1 + mm{d}_2 \ m{r}_i &= x_im{i} + y_im{j} \end{aligned}$$

Relative Position

 $m{r}_{ij} = m{r}_j - m{r}_i$

Generalized DOFs

$$\boldsymbol{q}_i = [x_i, y_i, \theta_i]^T$$



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Mechanical Interactions

• Kinetic energy:

$$T_i = \frac{1}{2}m(\dot{x}_i^2 + \dot{y}_i^2) + \frac{1}{2}I\dot{\theta}_i^2$$

• Potential energy

$$U_{ij}^{(e)} = \frac{1}{2}k_a(r_{ij} - r_{0_{ij}})^2 + \frac{1}{2}\alpha k_a \langle r_c - r_{ij} \rangle^2 + \frac{1}{2}k_\tau(\theta_i - \theta_b)^2 + \frac{1}{2}k_\tau(\theta_j - \theta_b)^2$$

$$\langle \bullet \rangle = \begin{cases} 0 & \bullet < 0 \\ \bullet & \ge 0 \end{cases} \quad \cos \theta_b = \frac{r_{0_{ij}} \cdot r_{ij}}{r_{0_{ij}}r_{ij}}$$

$$F = \begin{cases} k_a & f_a \\ f_c & f_c \\ f_c$$

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Magnetic Interactions

• Magnetic Work



Magnetic Force(*)

$$\begin{aligned} \boldsymbol{f}_{ij}^{(m)} &= -\frac{\mu_0}{4\pi} \boldsymbol{\nabla} \left(\frac{\boldsymbol{m}_i \cdot \boldsymbol{m}_j}{r_{ij}^3} - 3 \frac{(\boldsymbol{m}_i \cdot \boldsymbol{r}_{ij})(\boldsymbol{m}_j \cdot \boldsymbol{r}_{ij})}{r_{ij}^5} \right) \\ \boldsymbol{f}_{ij}^{(m)} &= \frac{3\mu_0}{4\pi} \frac{m_i m_j}{r_{ij}^4} \boldsymbol{n} = \frac{f^{(m)}}{r_{ij}^4} \boldsymbol{n} \end{aligned}$$

 $W_{ij}^{(m)} = -\int \boldsymbol{f}_{ij}^{(m)} \cdot d\boldsymbol{r}_{ij}$

Magnetic Potential

$$U_{ij}^{(m)} = -\int_{\infty}^{r_{ij}} \frac{f^{(m)}}{r_{ij}^4} \boldsymbol{n} \cdot d\boldsymbol{r}_{ij} = \frac{1}{3} \frac{f^{(m)}}{r_{ij}^3}$$

Radius of influence

$$r_{\infty} = \left(\frac{f^{(m)}}{3\varepsilon_e^{(m)}}\right)^{1/3}$$

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*An Analytic Solution for the Force Between Two Magnetic Dipoles, Yung, et al., 1998

Dissipative Interactions

• Dissipative Work

$$W_{ij}^{(d)} = -\int \boldsymbol{f}_{ij}^{(d)} \cdot d\boldsymbol{r}_{ij} - \int \phi_i^{(d)} d\theta_i - \int \phi_j^{(d)} d\theta_j$$

• Linear Damping

$$\boldsymbol{f}_{ij}^{(d)} = \begin{cases} c_a \dot{r}_{ij} \boldsymbol{\hat{n}} & r_{ij} > r_c \\ \beta c_a \dot{r}_{ij} \boldsymbol{\hat{n}} & r_{ij} \le r_c \end{cases}$$

• Angular Damping

$$\phi_i^{(d)} = -c_\tau (\dot{\theta}_i - \dot{\theta}_b)$$

Equations of Motion

• Total Energy

$$\mathcal{E} = \sum_{e=1}^{E} \mathcal{L}_e + W_e^{(d)}$$

• EOM for one particle

$$m{M}_i \ddot{m{q}}_i + \sum_{e_i}^{E_i} \left[m{C}_{e_i}(m{q}_{e_i}) \dot{m{q}}_{e_i} + m{K}_{e_i}(m{q}_{e_i}) m{q}_{e_i}
ight] = m{f}_i^{(m)}(m{q}_i, ..., m{q}_{m_i}, ..., m{q}_{M_i})$$

 Equation of motion for a particle is extended to a unit cell, and to a lattice

Equilibrium Identification

• Lattice Potential Energy

$$\mathcal{E} = \sum_{i=1}^{E} U_i^{(e)} + U_i^{(m)}$$

• Unit cell DOFs

$$oldsymbol{q} = [oldsymbol{q}_e,...,oldsymbol{q}_E]^T$$



• Periodic Lattice conditions are enforced on a specified unit cell size

$$oldsymbol{u} = [oldsymbol{q},oldsymbol{d}_1,oldsymbol{d}_2]$$

Optimization Problem

$$\min_{\boldsymbol{u}} \quad \mathcal{E}(\boldsymbol{u})$$

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Equilibrium Configurations



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2D Equilibrium Configuration Examples: 2 Particle Cells



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Energy profile of particle pair



Energy profile of particle pair, no contact





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Expansion through bistability





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Bistable Beam: shape control



Attractive Magnetic Forces



2D System Reconfiguration



Numerical Experiments

Hexagonal Lattice **Re-entrant Lattice** • •

Kagome lattice: model

• Lumped parameter model









Bloch Analysis

• Linearized magnetic forces

$$\bar{\boldsymbol{f}}^{(m)} = \boldsymbol{f}_0^{(m)} - 4\frac{f^{(m)}}{r_0^5}(r - r_0)\hat{\boldsymbol{n}} \qquad \theta_b = \frac{y'_j - y'_i}{r_{ij}}$$

• EOM for a Unit Cell

$$M\ddot{q}_{n,m} + \sum_{n=-N}^{N} \sum_{m=-M}^{M} \left(K_{n,m}^{(e)} + K_{n,m}^{(m)}
ight) q_{n,m} = 0$$

Plane wave solution

$$\boldsymbol{q}_{n,m} = \boldsymbol{q}_0 e^{i(\omega t + \boldsymbol{\kappa} \cdot (n\boldsymbol{d}_1 + m\boldsymbol{d}_2))}$$

• Linear EVP $\begin{bmatrix} -\omega^2 M + \sum_{n,m} \left(K_{n,m} e^{i\boldsymbol{\kappa} \cdot (n\boldsymbol{d}_1 + m\boldsymbol{d}_2)} \right) \end{bmatrix} \boldsymbol{q}_0 = \boldsymbol{0}$ Non-local interactions $\begin{pmatrix} \boldsymbol{K}(\boldsymbol{\kappa}) - \omega^2 \boldsymbol{M} \end{pmatrix} \boldsymbol{q}_0 = \boldsymbol{0} \qquad \boldsymbol{K}(\boldsymbol{\kappa}) = \sum_{n,m} \left(\boldsymbol{K}_{n,m} e^{i\boldsymbol{\kappa} \cdot (n\boldsymbol{d}_1 + m\boldsymbol{d}_2)} \right)$

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Effect of Reconfiguration on Band Structure



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Wave Propagation in 3D Structures



Krödel, S., Delpero, T., Bergamini, A., Ermanni, P., & Kochmann, D. M. (2013). 3D Auxetic Microlattices with Independently Controllable Acoustic Band Gaps and Quasi-Static Elastic Moduli. Advanced Engineering Materials, 15(9999).

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Wave Propagation in 3D Lattices



Wave Propagation in 3D Lattices



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Wave Propagation in 3D Lattices



Effect of morphology on wave motion Flat plate **Curved plate**





- Band diagram, flat plate
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.01$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.02$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.03$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.04$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.05$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.06$
 - Black line: 100 elements
 - Blue line: 2500 elements





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- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.07$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.08$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.09$
 - Black line: 100 elements
 - Blue line: 2500 elements





- Band diagram, $\alpha = 0^{\circ}$, $\eta = 0.10$
 - Black line: 100 elements
 - Blue line: 2500 elements





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Summary

- Adaptive wave properties are highly desirable in a variety of applications
- Adaptivity can be achieved from controlled topological changes associated with:
 - large deformations and
 - multifield interactions
- Magneto-elastic lattices provide a good framework for these studies
- Can a "material" concept be derived from such configurations?



Thank you



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